

# NANOPHYSIQUE

## INTRODUCTION PHYSIQUE AUX NANOSCIENCES

### *2. PRINCIPALES METHODES DE MICROSCOPIE*

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# *METHODES DE MICROSCOPIE*

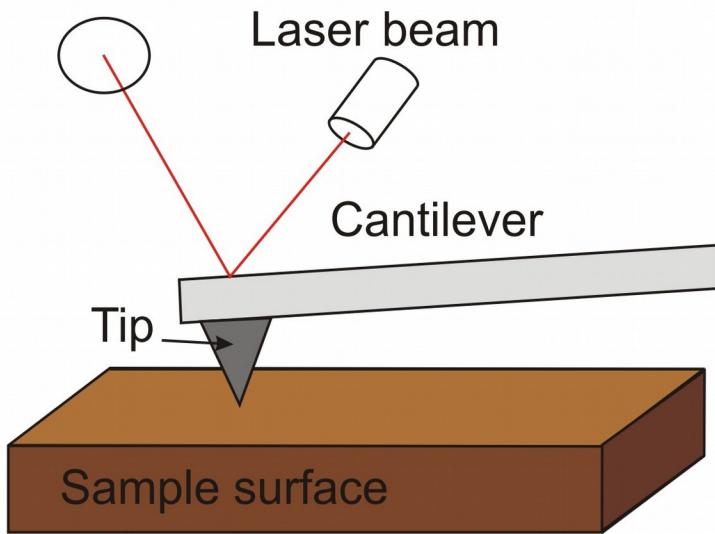
- **Paramètres Fondamentaux**
- **Microscopes Optiques**
  - **Principe**
  - **Améliorations:** phase contrast, dark field, fluorescent, ...
  - **Cristallographe aux Rayon X**
- **Microscope Electronique**
  - **à Transmission**
  - **à Balayage**
- **Microscope à émission champ**
- **Microscope à effet tunnel électronique**
- **Microscope à force atomique**
- **Optical Tweezers**
- **Light Scattering**

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# Microscope à force atomique

Photodetector



Laser beam

Cantilever

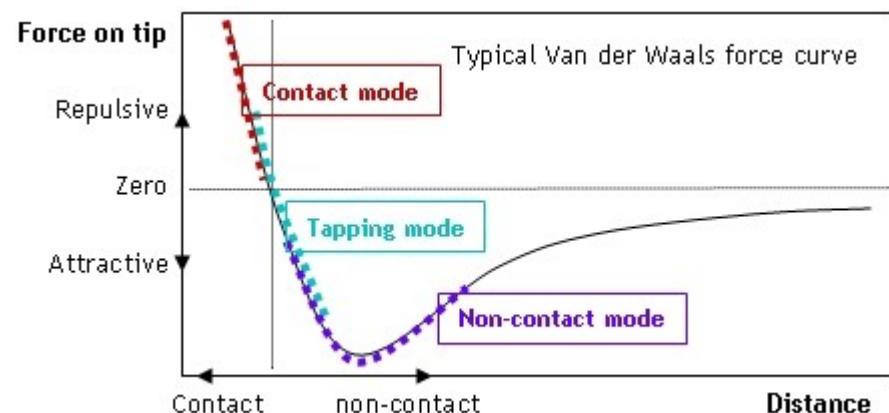
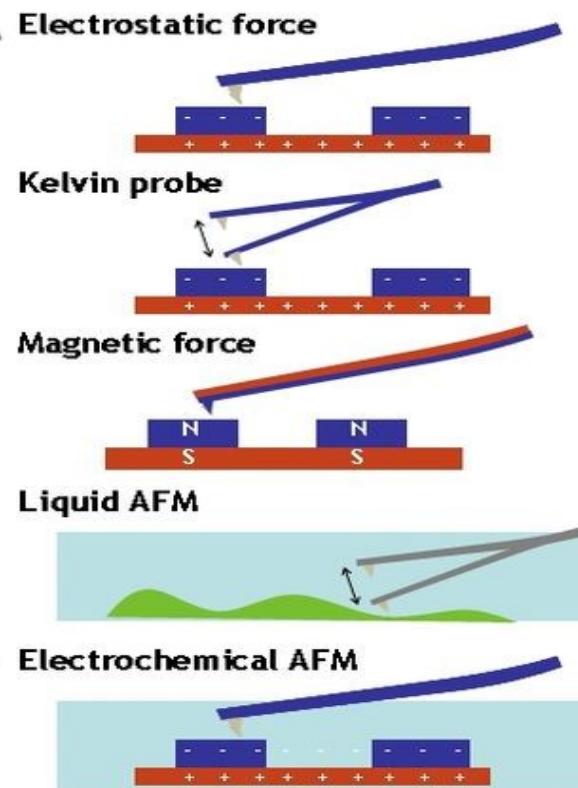
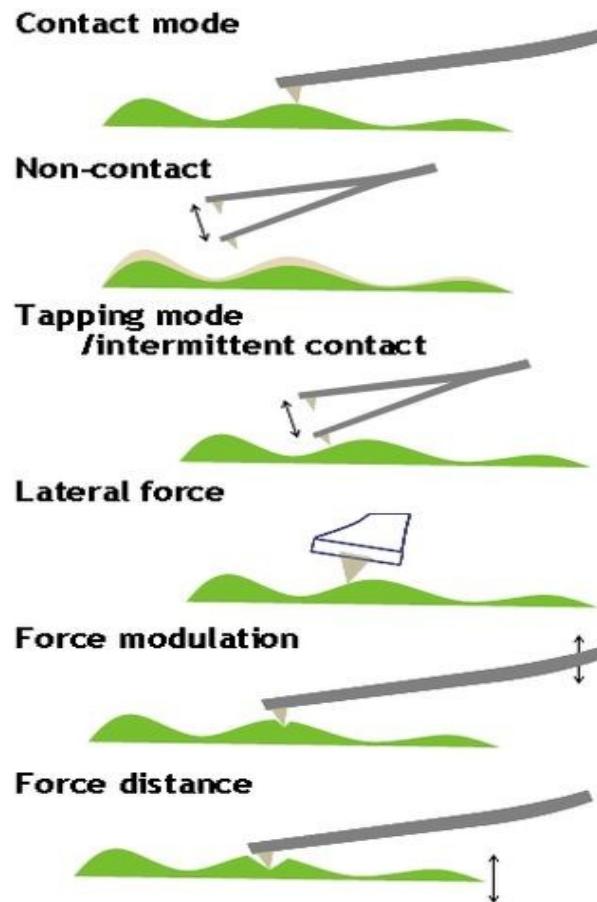
Tip

Sample surface

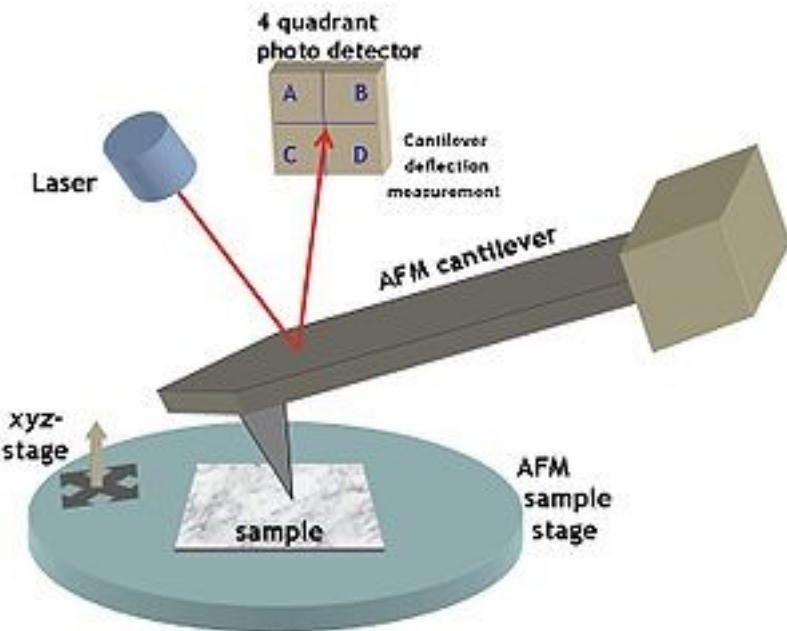
|                     | <b>STM</b>         | <b>AFM</b>              |
|---------------------|--------------------|-------------------------|
| Lateral Resolution  | 0.5-1 nm           | 0.5 nm                  |
| Vertical Resolution | 2D only            | 0.05nm                  |
| Field of view       | 1-2 X 1-2 mm       | 100 x 100 $\mu\text{m}$ |
| Vertical range      |                    | 100 $\mu\text{m}$       |
| Preparation         | Couche conductrice | ----                    |
| Environment         | vide               | L'air, liquide          |

Kurganskaya, I.; Luttge, A.; Barron, A. The Application of VSI (Vertical Scanning Interferometry) to the Study of Crystal Surface Processes, Connexions Web site. <http://cnx.org/content/m22326/1.4/>, Jul 13, 2009.

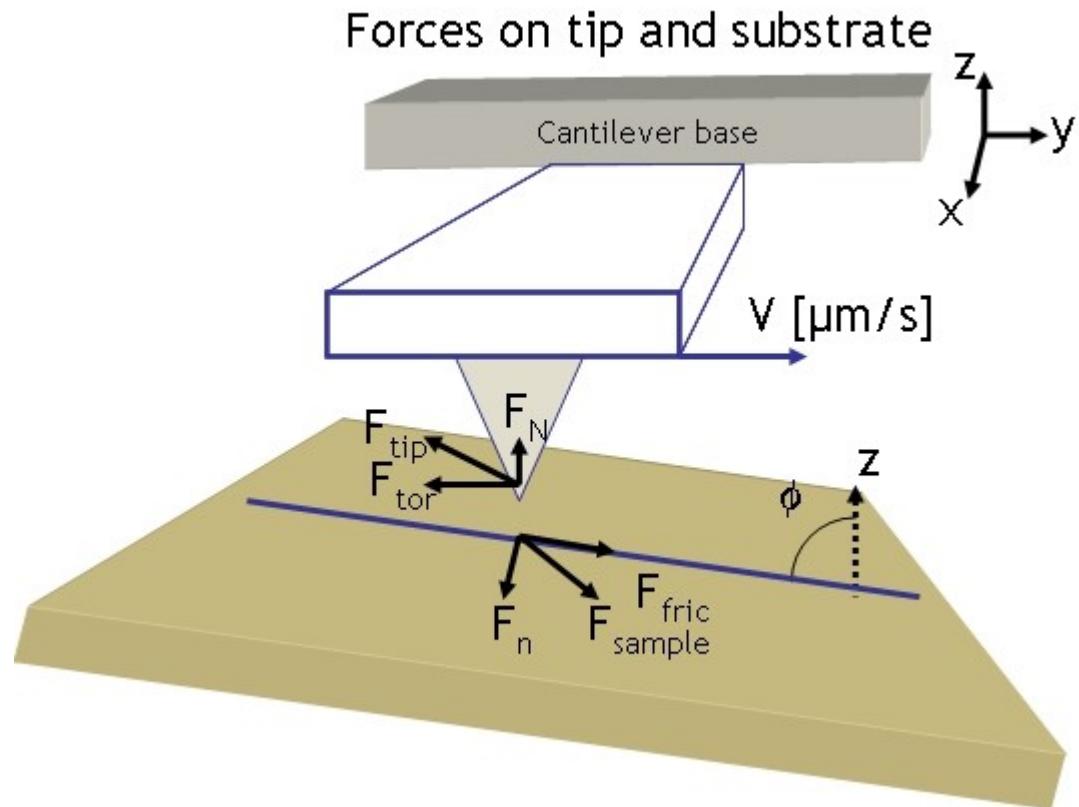
# AFM: les modes de fonctionnement



# Microscope à force atomique



Signal de droit-gauche:  $A+C-(B+D)$   
Signal de haut en bas:  $A+B-(C+D)$



# Microscope à force atomique

## I. PRINCIPE GÉNÉRAL: UNE OSCILLATEUR CLASSIQUE

$$\ddot{u} + 2\beta\dot{u} + \omega_0^2 u = \gamma \cos \omega t + \frac{1}{m} F(D, u)$$

où

$D$  = distance entre la surface et la position de la pointe quand le cantilever n'est pas défléchi.

$z$  = distance entre la surface et la position de la pointe actuelle

$u = z - D$  = déviation

$m$  = mass effictive

$\omega_0 = \sqrt{\frac{k}{m}}$  = la fréquence de résonanance de l'oscillateur

$k$  = la raideur du cantilever

$\beta$  = un terme de dissipation

$\gamma$  = l'amplitude de l'excitation

$\omega$  = fréquence de l'excitation

$F(D, u)$  = la force d'ineteraction pointe-surface

N.B.  $Q \equiv \frac{\omega_0}{2\beta}$  est *le facteur de qualité*.

# Microscope à force atomique

## II. CASE I: CONTACT MODE

Ne pas d'excitation:

$$\ddot{u} + 2\beta\dot{u} + \omega_0^2 u = \frac{1}{m} F(D, u) \implies ku = F(D, u)$$

e.g.

$$ku \simeq F(D) + uF'(D) \implies u = \frac{F(D)}{k - F'(D)}, \text{ Si } k \gg F'(D), u \simeq \frac{F(D)}{k}$$

## III. CASE II: LE MODE RÉSONNANT LINÉAIRE

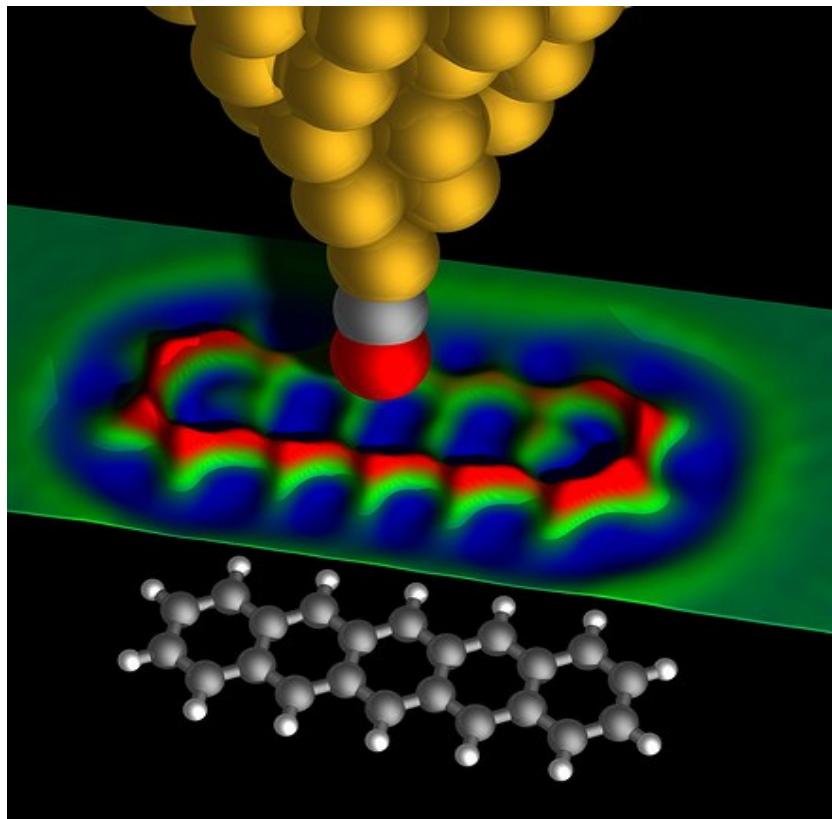
$$\ddot{u} + 2\beta\dot{u} + \omega_0^2 u \simeq \gamma \cos \omega t + \frac{1}{m} F(D) + u \frac{1}{m} F'(D)$$

de sorte que

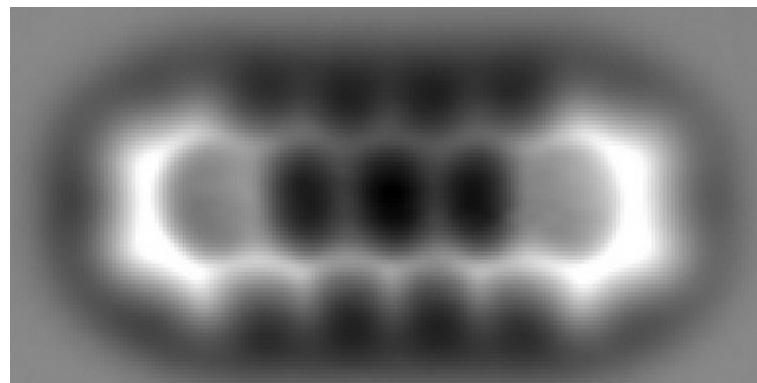
$$\ddot{u} + 2\beta\dot{u} + \omega_0^2 \left(1 - \frac{1}{k} F'(D)\right) u \simeq \gamma \cos \omega t + \frac{1}{m} F(D)$$

Ça veux dire que il y a un changement de fréquence naturelle de l'oscillateur.

# AFM Recherche actuelle ...



Imaging the "anatomy" of a pentacene molecule - 3D rendered view: By using an atomically sharp metal tip terminated with a carbon monoxide (CO) molecule, IBM scientists were able to measure in the short-range regime of forces which allowed them to obtain an image of the inner structure of the molecule. The colored surface represents experimental data. (Image courtesy of IBM Research/Zurich)



# Resume

|                     | <b>Optique</b>      | <b>Xray</b>    | <b>Confocal</b>  | <b>TEM/SEM</b> | <b>STM</b>         | <b>AFM</b>     |
|---------------------|---------------------|----------------|--|----------------|--------------------|----------------|
| Lateral Resolution  | 200nm               | 25nm           | 200nm  | 0.1nm/3nm      | 0.1 nm             | 0.5 nm         |
| Vertical Resolution | 2D only             | -----          | 500nm  | -----          | 2D only            | 0.05nm         |
| Field of view       | grande              | 50µm           | grande   | Baylage        | 1-2 X 1-2 mm       | 100 x 100 µm   |
| Vertical range      | -----               | -----          | Limité par le temps<br>(1-1000 sec/mm <sup>2</sup> /tranche) | -----          | -----              | 100 µm         |
| Preparation         | -----               | -----          | -----  | tres mince     | Couche conductrice | ---            |
| Environment         | L'air, liquide, ... | L'air, liquide | liquide  | vide           | vide               | L'air, liquide |

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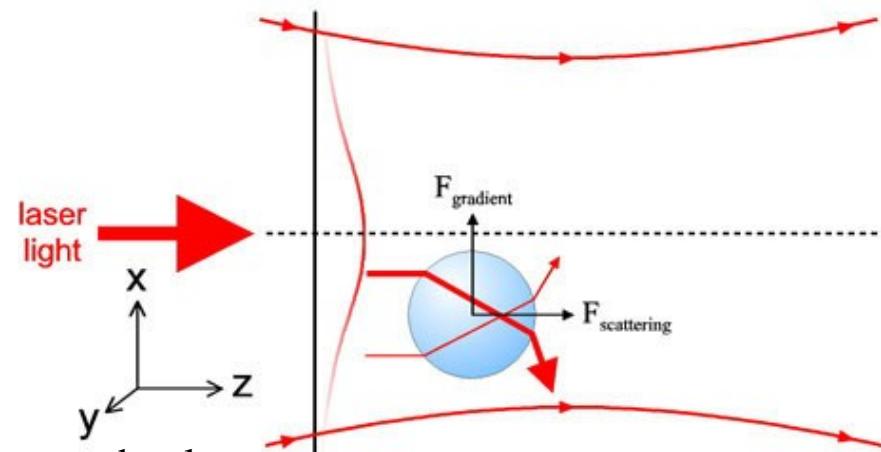
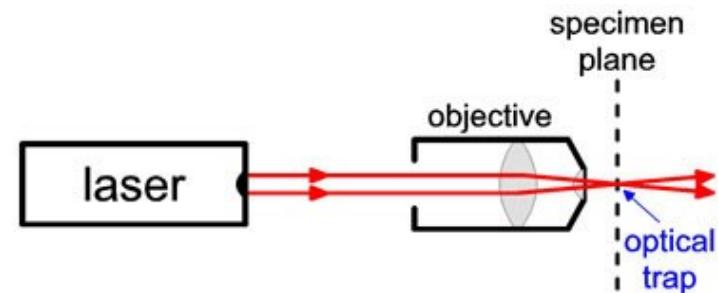
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# Optical Tweezers

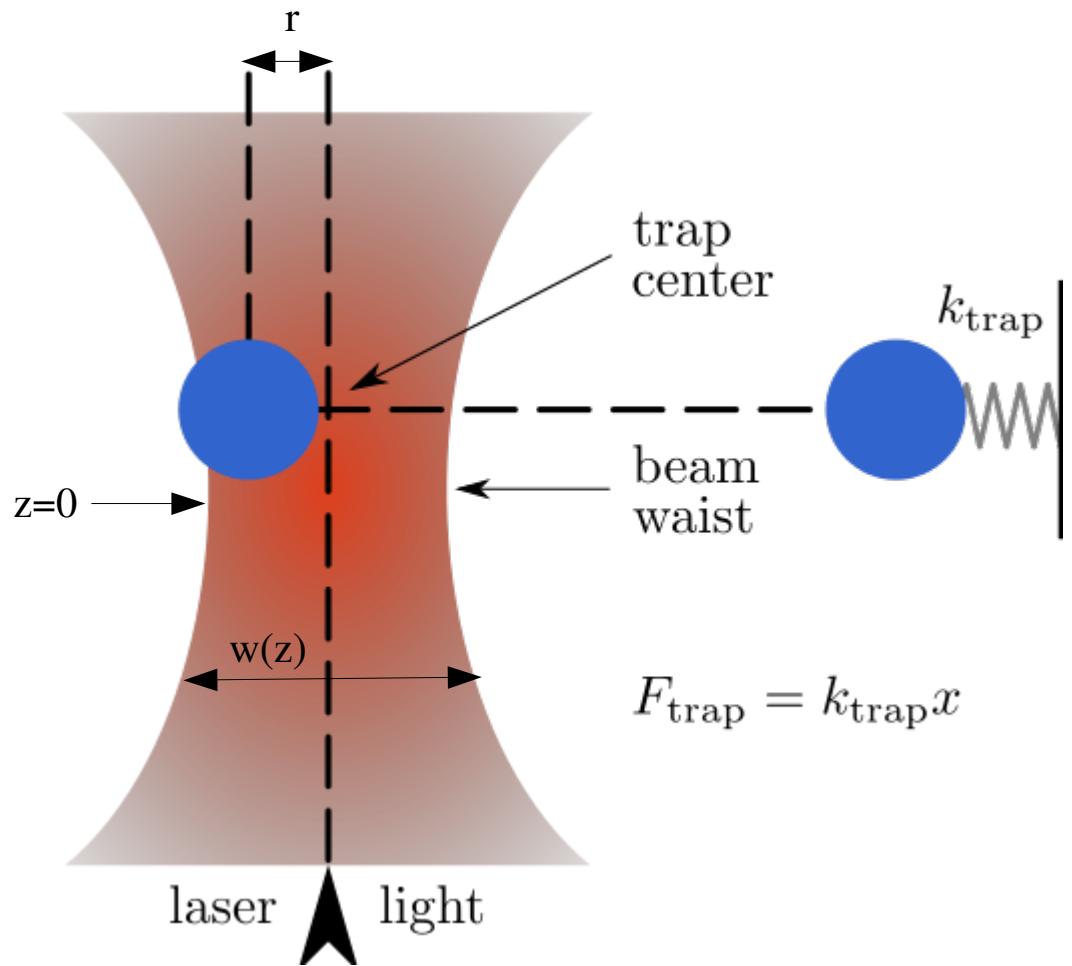
**Optical Tweezers** use light to manipulate microscopic objects as small as a single atom. The radiation pressure from a focused laser beam is able to trap small particles. In the biological sciences, these instruments have been used to apply forces in the pN-range and to measure displacements in the nm range of objects ranging in size from 10 nm to over 100 mm.

3 regimes:

- $D \gg \lambda \implies$  ray optics
- $D \sim \lambda \implies$  Maxwell's equations
- $D \ll \lambda \implies$  Electrostatics



# Optical Tweezers

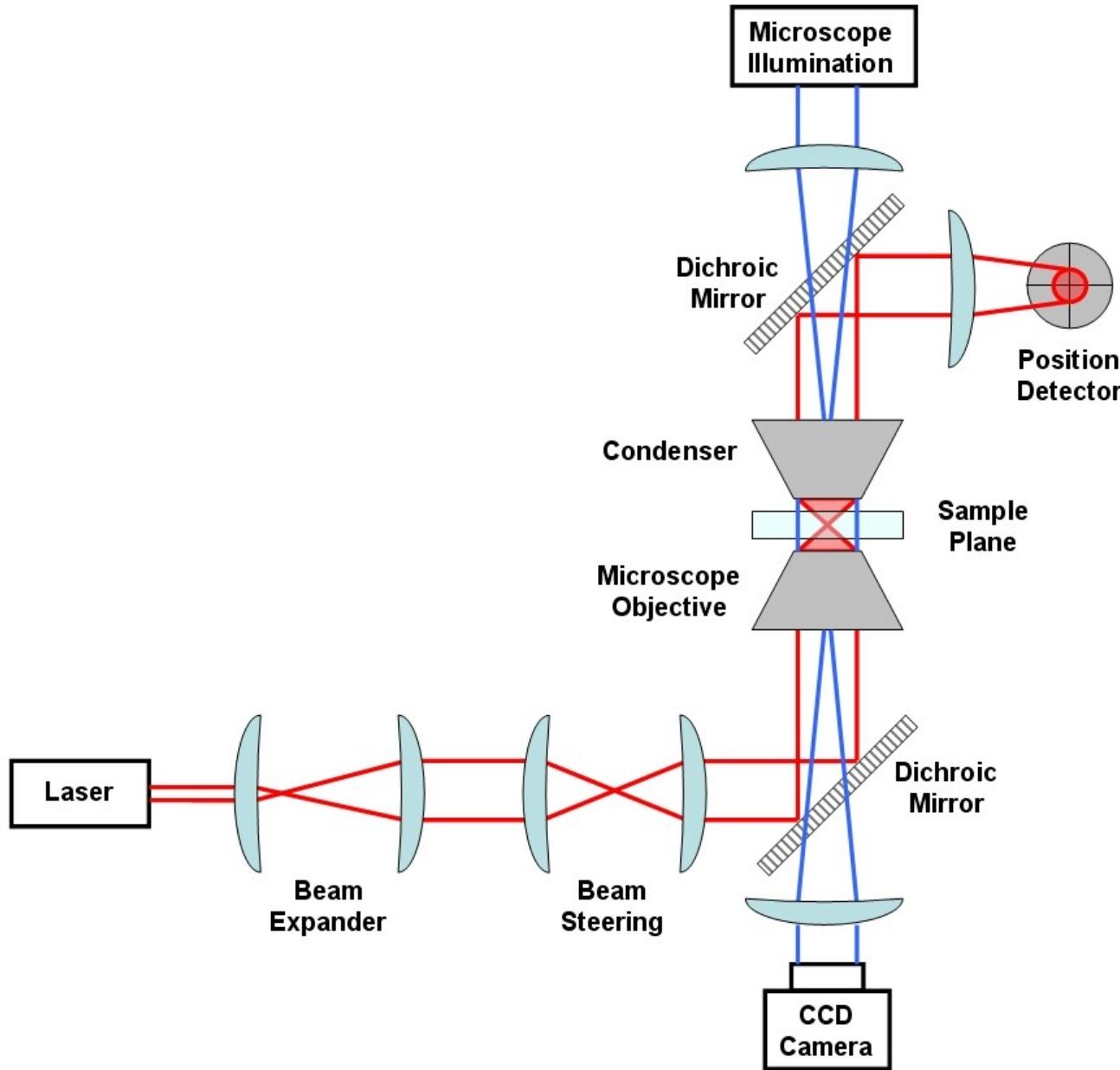


$$I(\mathbf{r}) = \frac{1}{2\eta} E^2(\mathbf{r}) = I_0 \left( \frac{w_0}{w(z)} \right)^2 \exp \left( -\frac{2r^2}{w(z)^2} \right)$$

$$w(z) = w_0 \sqrt{1 + \left( \frac{z}{z_R} \right)^2}$$

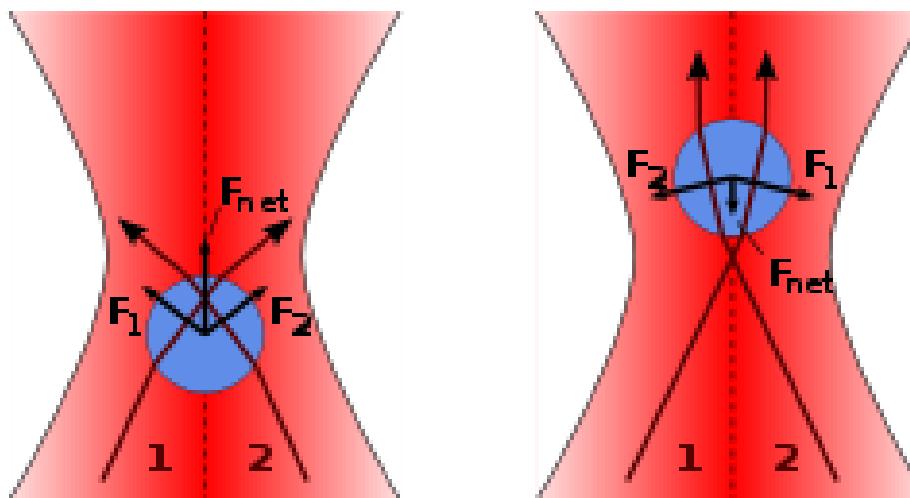
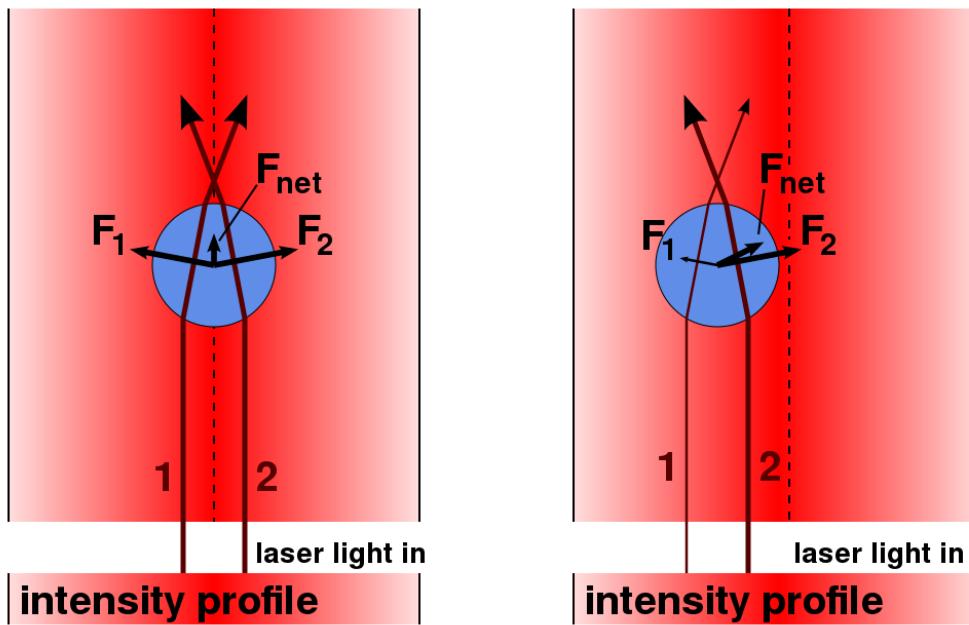
$$z_R = \frac{\pi w_0^2}{\lambda}$$

# Optical Tweezers



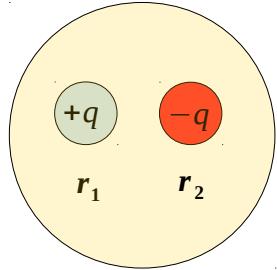
# Optical Tweezers: $D \gg \lambda \implies$ ray optics

Force is due to change in momentum of refracted light.



# Optical Tweezers: $D \ll \lambda$

Particle is treated as a point (induced) dipole



$$\mathbf{R} = \frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \quad \mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2 \quad \mathbf{d} = q \mathbf{r}$$

$$\mathbf{F}_i = q_i \left\{ \mathbf{E}(\mathbf{r}_i) + \frac{d \mathbf{r}_i}{dt} \times \mathbf{B}(\mathbf{r}_i) \right\}$$

$$\mathbf{F}_{total} = q \left\{ \mathbf{r} \cdot \nabla \mathbf{E}(\mathbf{R}) + \frac{d \mathbf{r}}{dt} \times \mathbf{B}(\mathbf{R}) \right\} + \text{higher order in } r$$

$$= \mathbf{d} \cdot \nabla \mathbf{E}(\mathbf{R}) + \frac{d \mathbf{d}}{dt} \times \mathbf{B}(\mathbf{R}) + \text{higher order in } r$$

Assuming linear dielectric:  $\mathbf{d} = \alpha \mathbf{E}$

and using one of Maxwell's equations:  $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$

$$\mathbf{F}_{total} = \alpha \left\{ \nabla \mathbf{E}(\mathbf{R})^2 + \underbrace{\frac{\partial}{\partial t} (\mathbf{E}(\mathbf{R}) \times \mathbf{B}(\mathbf{R}))}_{\text{power per unit area}} \right\} + \text{higher order in } r$$

For dielectric sphere

$$\alpha = \frac{\pi D^3 \epsilon_0}{2} \frac{\epsilon - \epsilon_0}{\epsilon + 2 \epsilon_0}$$

$$\approx \frac{\pi D^3 \epsilon_0}{2} \frac{n^2 - n_0^2}{n^2 + 2n_0^2}$$

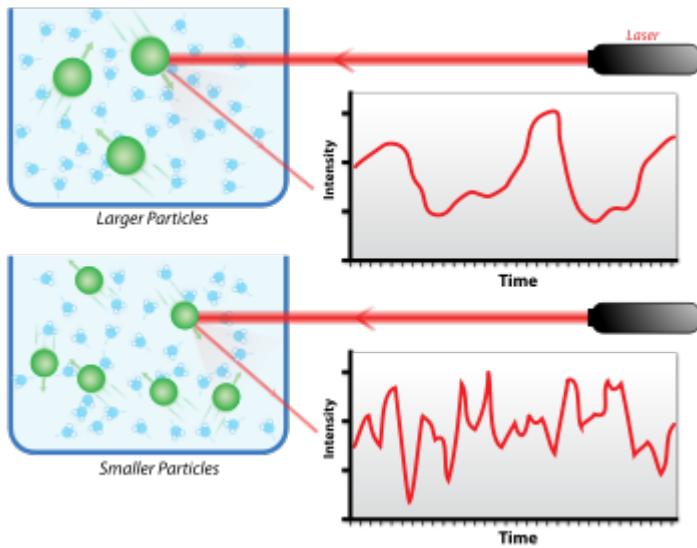
Proof = exercise!

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# Dynamic light scattering

- Typically used for particles diffusing in a liquid bath
- Determines size of particles



Field auto-correlation function (what you want):

$$g_1(q,t) = \frac{\langle E(q,t)E(q,t+\tau) \rangle}{\langle E(q,t)E(q,t) \rangle}$$

Intensity auto-correlation function (what you measure):

$$g_2(q,t) = \frac{\langle I(q,t)I(q,t+\tau) \rangle}{\langle I(q,t)I(q,t) \rangle}$$

$$g_2(q,t) \sim 1 + \text{const} \times [g_1(q,t)]^2$$

$$g_1(q,t) = \exp(-q^2 D t), \quad D = \text{diffusion constant}$$

$$q = \frac{4\pi n_0}{\lambda} \sin\left(\frac{\theta}{2}\right)$$